

Generalized parton distributions from
domain wall valence quarks and staggered sea quarks

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Mixed Action Lattice Calculation

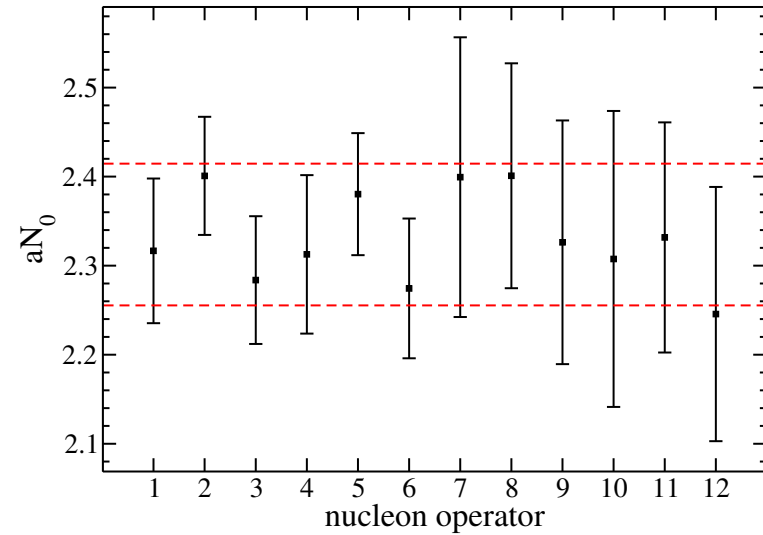
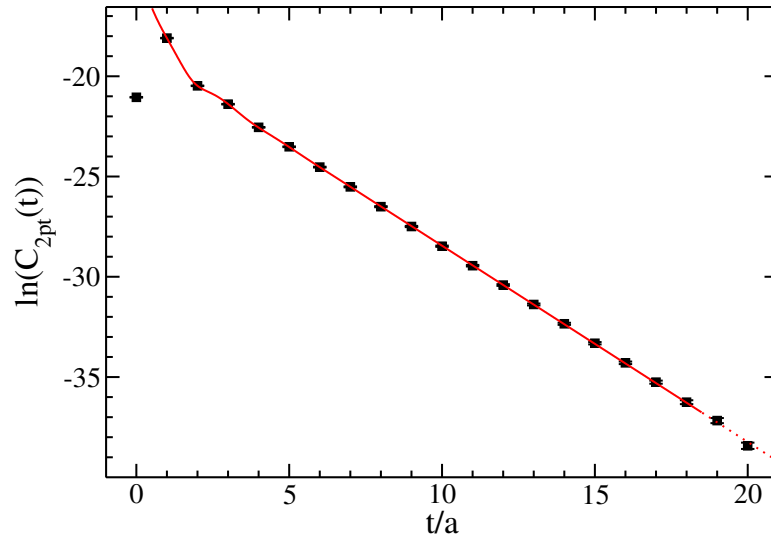
- asqtad staggered sea quarks (MILC) with $a = 0.124$ fm

$am_{u/d}^{\text{asqtad}}$	L/a	L	m_{π}^{DWF}	#
		fm	MeV	
0.05	20	2.52	761	425
0.04	"	"	693	350
0.03	"	"	594	564
0.02	"	"	498	486
0.01	"	"	354	656
0.01	28	3.53	353	270

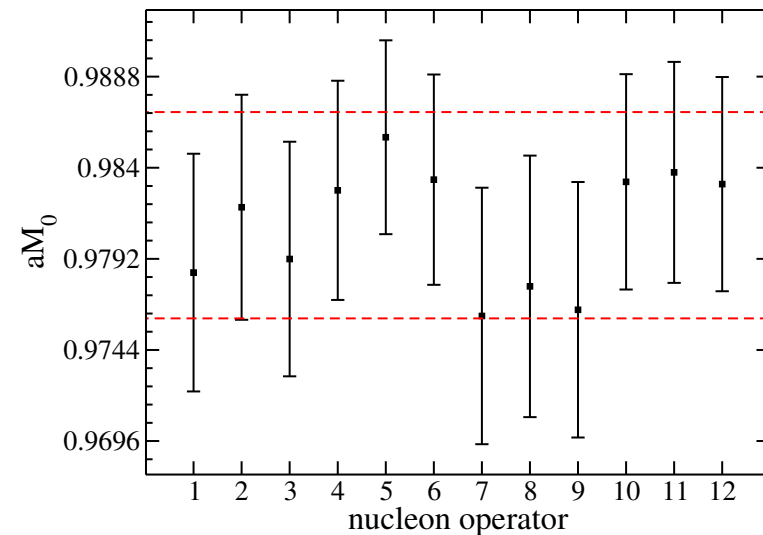
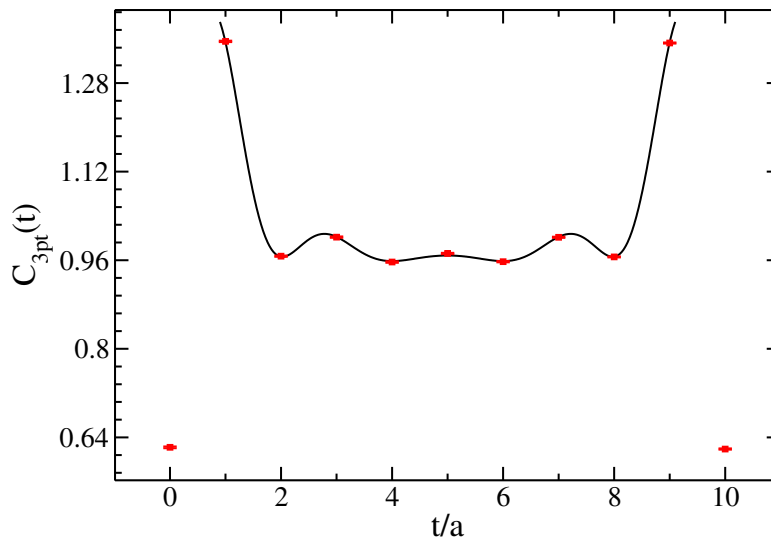
- domain wall valence quarks with HYP smearing and $L_5 = 16$
- one-loop perturbative renormalization at $\mu = 2$ GeV

Oscillatory Domain Wall Correlators and Fits

- $$C_{2\text{pt}}(t) = A_0 e^{-M_0 t} + A_1 e^{-M_1 t} + B_0 (-1)^t e^{-N_0 t}$$



- $$C_{3\text{pt}}(t) = O_{00} + O_{10} \cosh((M_1 - M_0)t) + P_{00} (-1)^t \cosh((N_0 - M_0)t)$$



Moments of Parton Distributions

- forward matrix elements of twist two operators

$$O_q^{\mu_1 \cdots \mu_n} = \bar{q} i D^{(\mu_1} \cdots i D^{\mu_{n-1}} \gamma^{\mu_n)} q$$

$$\langle P, S | O_q^{\mu_1 \cdots \mu_n} | P, S \rangle = 2 \langle x^{n-1} \rangle_q P_{\mu_1} \cdots P_{\mu_n}$$

- unpolarized, helicity and transversity moments

$$\langle x^n \rangle_q = \int_{-1}^1 dx x^n q(x)$$

$$\langle x^n \rangle_{\Delta q} = \int_{-1}^1 dx x^n \Delta q(x)$$

$$\langle x^n \rangle_{\delta q} = \int_{-1}^1 dx x^n \delta q(x)$$

Axial Charge and Chiral Perturbation Theory

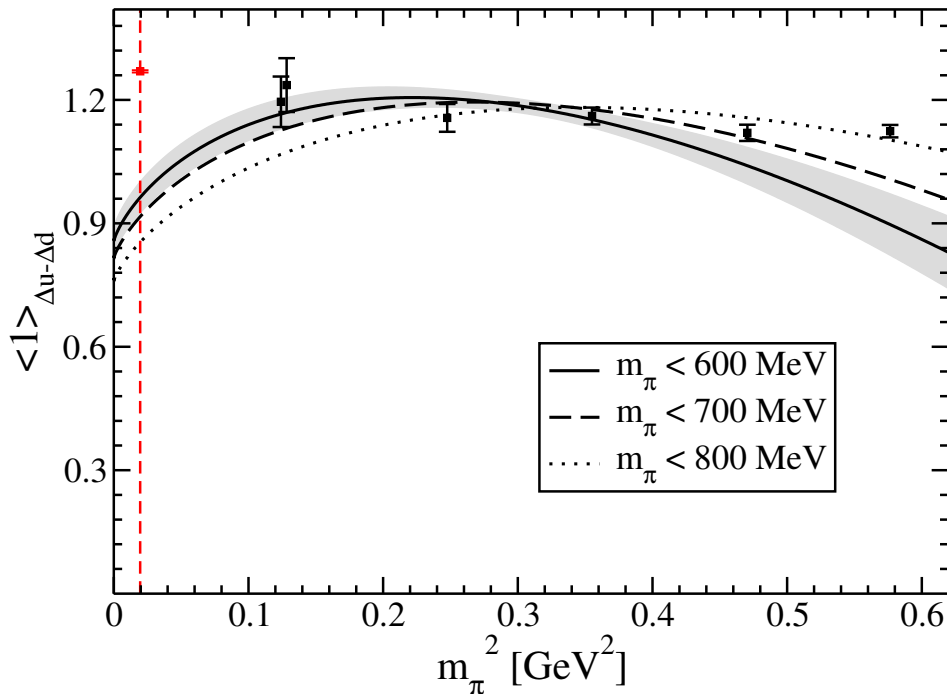
- standard one-loop heavy baryon chipt

$$g = \hat{g} \left(1 - \frac{(2\hat{g}^2 + 1)}{(4\pi\hat{f})^2} m^2 \ln \left(\frac{m^2}{\mu^2} \right) \right) + c_g(\mu) m^2$$

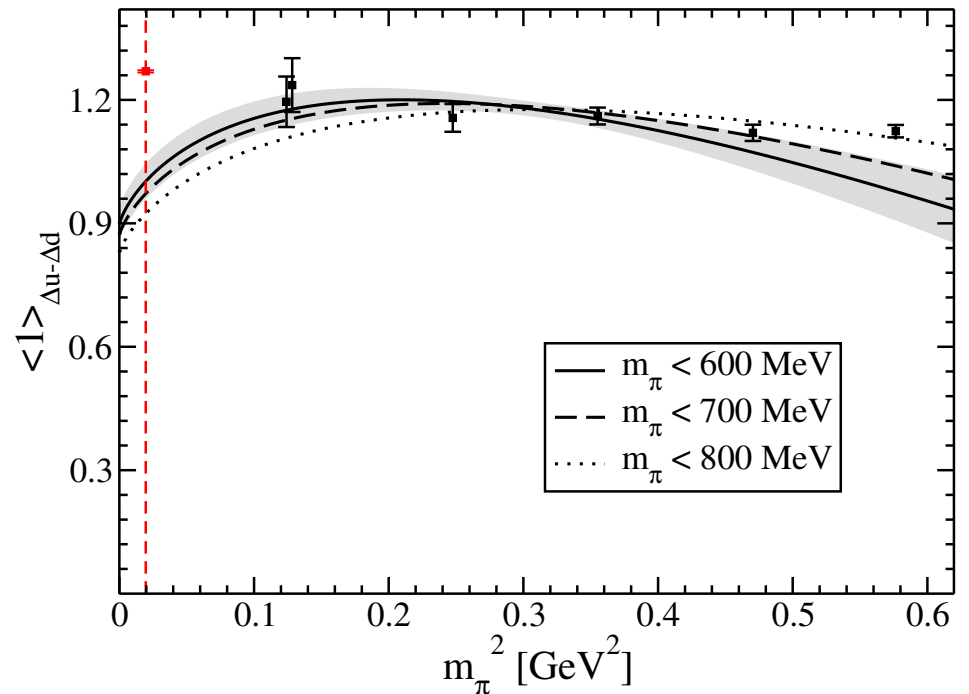
- finite range regulated one-loop heavy baryon chipt

$$g = \hat{g} \left(1 - \frac{(2\hat{g}^2 + 1)}{(4\pi\hat{f})^2} m^2 \ln \left(\frac{m^2}{\Lambda^2 + m^2} \right) \right) + c_g(\Lambda) m^2$$

Heavy Baryon χ PT



Finite Range Regulator χ PT



Self Consistent Chiral Perturbation Theory

- choose $\mu = \alpha \overset{\circ}{f}$ and replace $\overset{\circ}{f}$, $\overset{\circ}{g}$, $\langle \overset{\circ}{x} \rangle$ with f , g , $\langle x \rangle$ in the NLO term

$$\langle x \rangle = \langle \overset{\circ}{x} \rangle - \langle \overset{\circ}{x} \rangle (3\overset{\circ}{g}^2 + 1)m^2 / (4\pi \overset{\circ}{f})^2 \ln(m^2/\mu^2) + c_x(\mu)m^2$$

$$\langle x \rangle (1 + (3g^2 + 1)m^2 / (4\pi f)^2 \ln(m^2/(\alpha f)^2)) = \langle \overset{\circ}{x} \rangle + c_x(\alpha)m^2$$

- choose $\mu = \alpha \overset{\circ}{f}$ and replace $\overset{\circ}{f}$, $\overset{\circ}{g}$ with f , g in the NLO term

$$g = \overset{\circ}{g} - (2\overset{\circ}{g}^3 + \overset{\circ}{g})m^2 / (4\pi \overset{\circ}{f})^2 \ln(m^2/\mu^2) + c_g(\mu)m^2$$

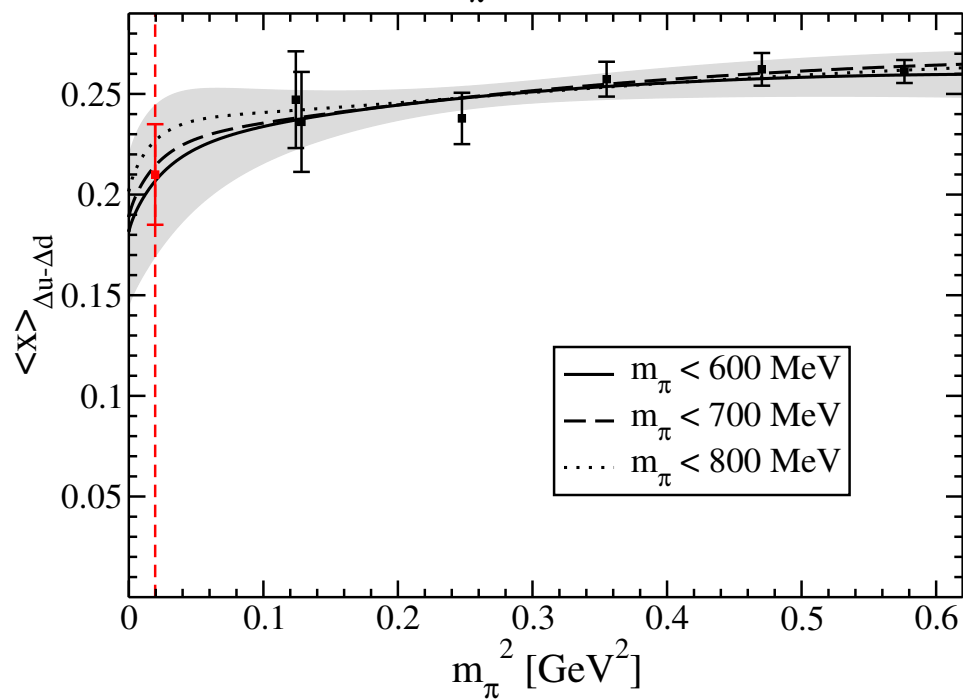
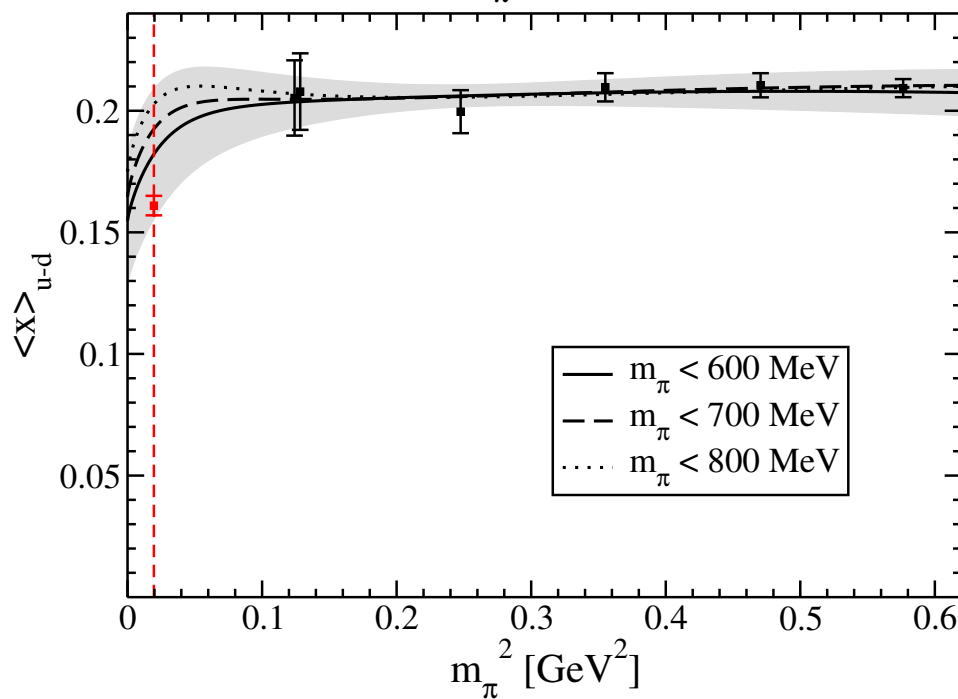
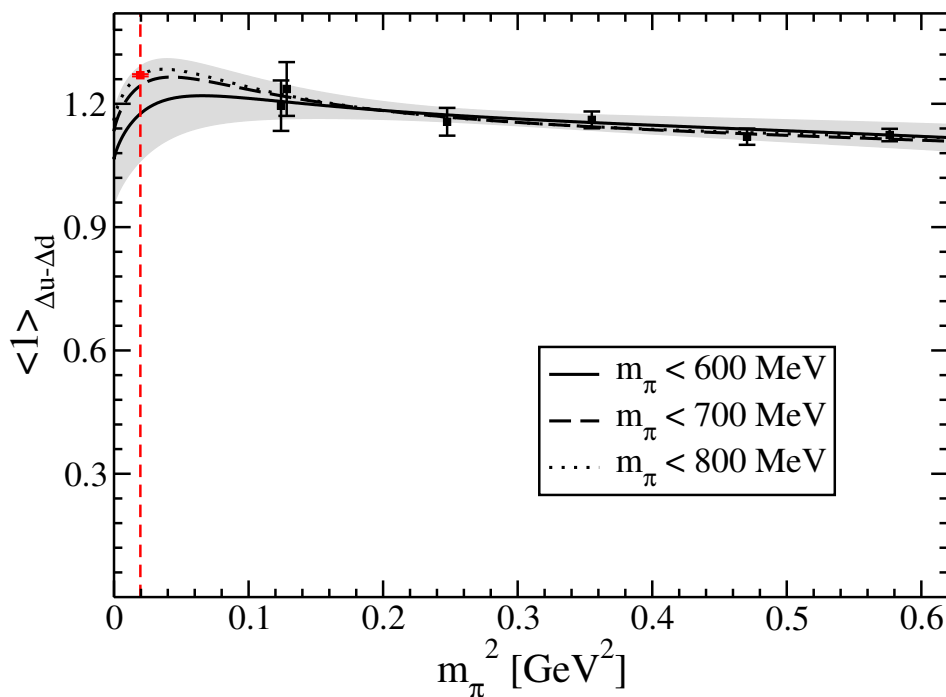
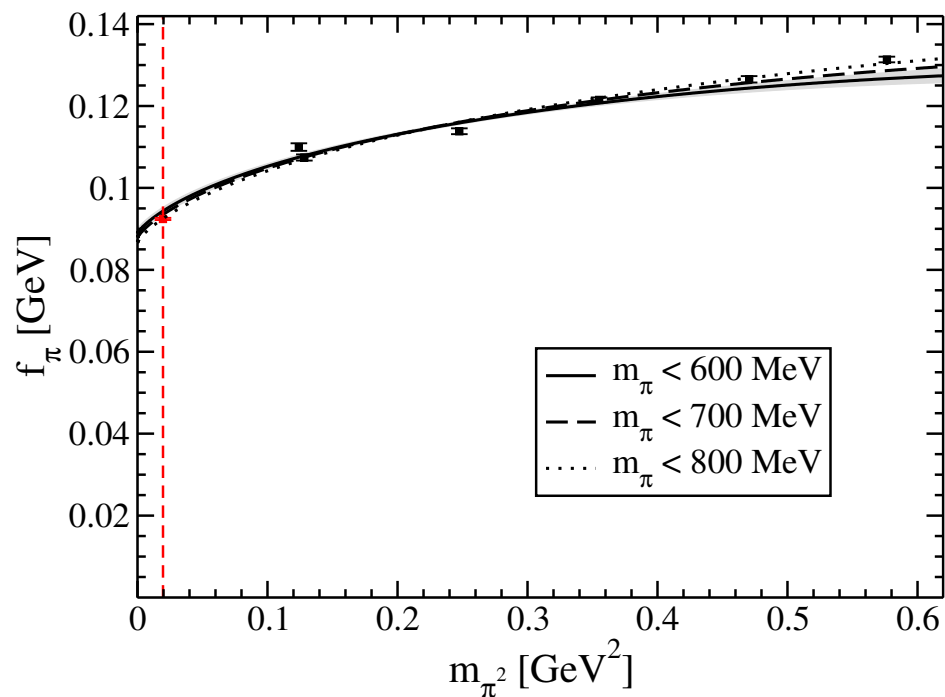
$$g(1 + (2g^2 + 1)m^2 / (4\pi f)^2 \ln(m^2/(\alpha f)^2)) = \overset{\circ}{g} + c_g(\alpha)m^2$$

- choose $\mu = \alpha \overset{\circ}{f}$ and replace $\overset{\circ}{f}$ with f in the NLO term

$$f = \overset{\circ}{f} - m^2 / (4\pi \overset{\circ}{f}) \ln(m^2/\mu^2) + c_f(\mu)m^2$$

$$f(1 + m^2 / (4\pi f)^2 \ln(m^2/(\alpha f)^2)) = \overset{\circ}{f} + c_f(\alpha)m^2$$

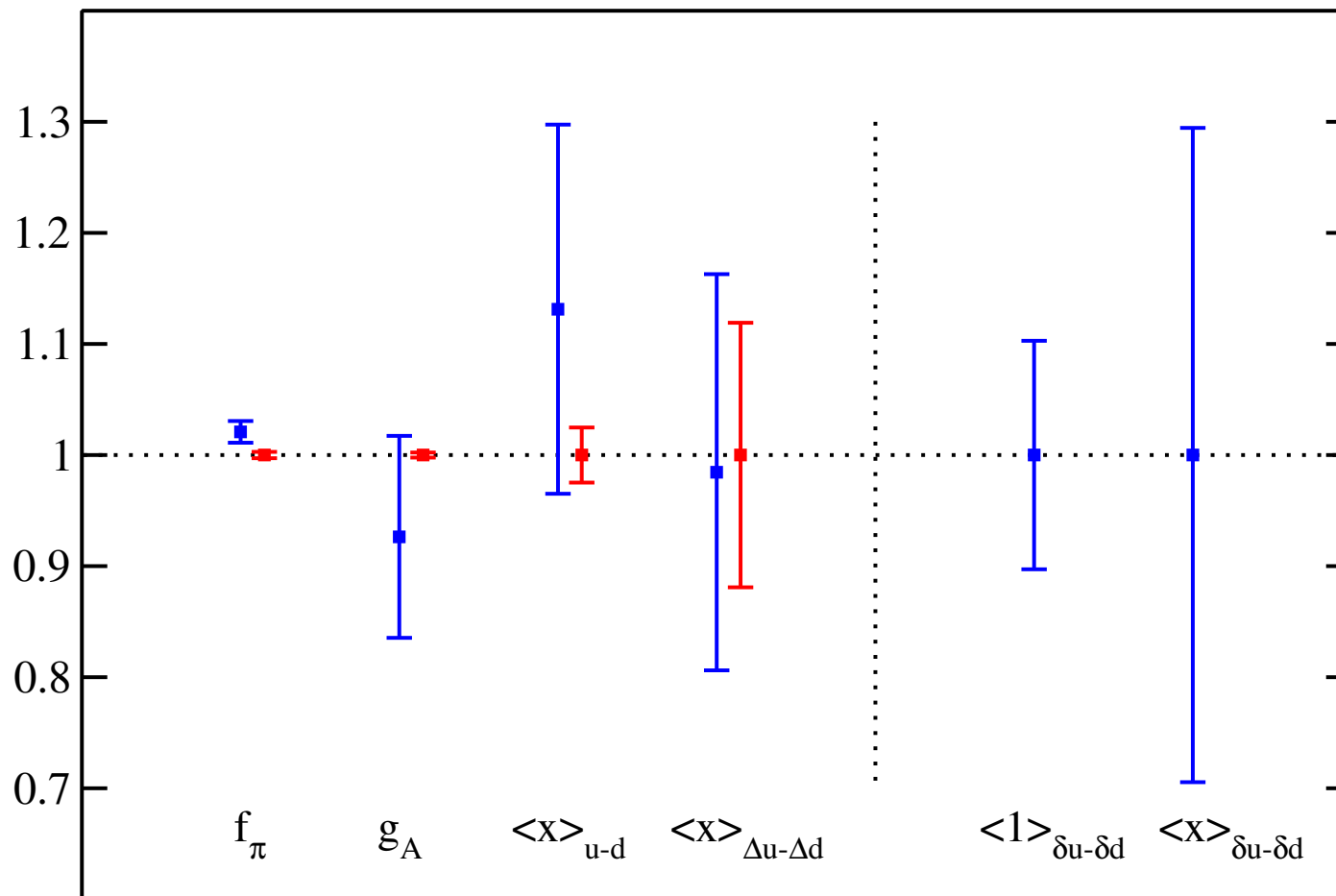
Self Consistent Chiral Perturbation Theory



Self Consistent Chiral Perturbation Theory

- self consistent χ PT successfully extrapolates f_π , g_A , $\langle x \rangle_{u-d}$, $\langle x \rangle_{\Delta u-\Delta d}$

Lattice / Experiment



- and highlights the potential for genuine predictions of $\langle 1 \rangle_{\delta u-\delta d}$, $\langle x \rangle_{\delta u-\delta d}$

Moments of Generalized Parton Distributions

- off-forward matrix elements of twist two operators

$$\langle P' | O_q^{\mu_1 \dots \mu_n} | P \rangle = \sum_{\substack{i=0 \\ \text{even}}}^{n-1} \left\{ A_{ni}^q(t) K_{ni}^A + B_{ni}^q(t) K_{ni}^B \right\} + \delta_{\text{even}}^n C_n^q(t) K_n^C$$

- quark angular momenta

$$\frac{1}{2} = \frac{1}{2} \Delta \Sigma^{u+d} + L^{u+d} + J^g$$

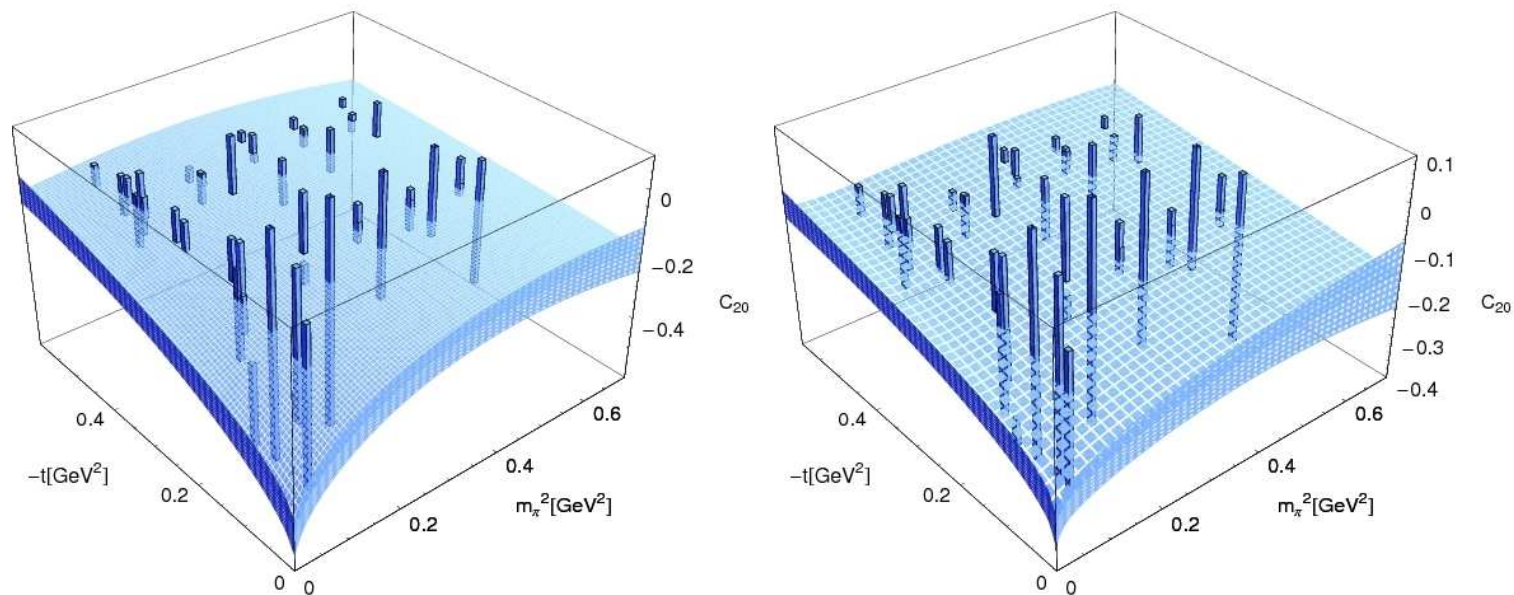
$$\Delta \Sigma^q = \tilde{A}_{10}^q(0) \quad J^q = \frac{1}{2} \left(A_{20}^q(0) + B_{20}^q(0) \right) \quad L^q = J^q - \frac{1}{2} \Delta \Sigma^q$$

- transverse quark distributions

$$\int_{-1}^1 dx x^{n-1} q(x, \vec{b}_\perp) = \int \frac{d^2 \Delta_\perp}{(2\pi)^2} e^{-i\vec{b}_\perp \cdot \vec{\Delta}_\perp} A_{n0}^q(-\vec{\Delta}_\perp^2)$$

Covariant versus Heavy Baryon Chiral Perturbation Theory

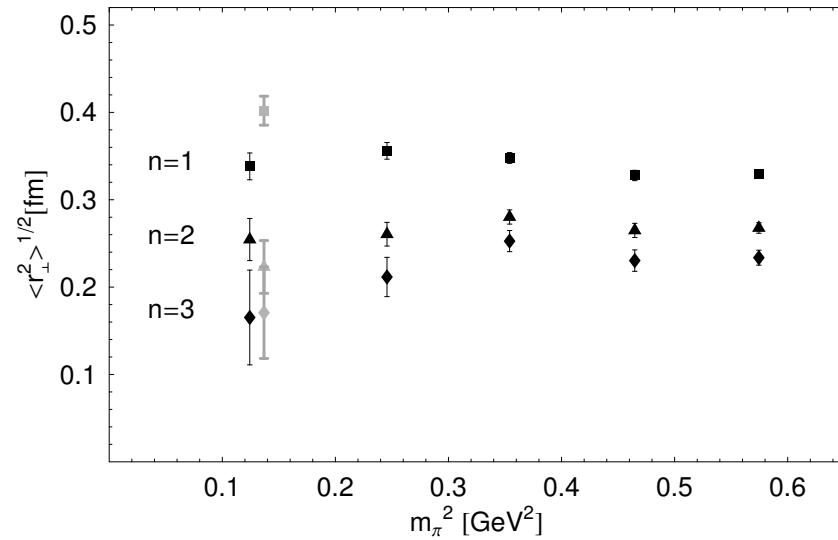
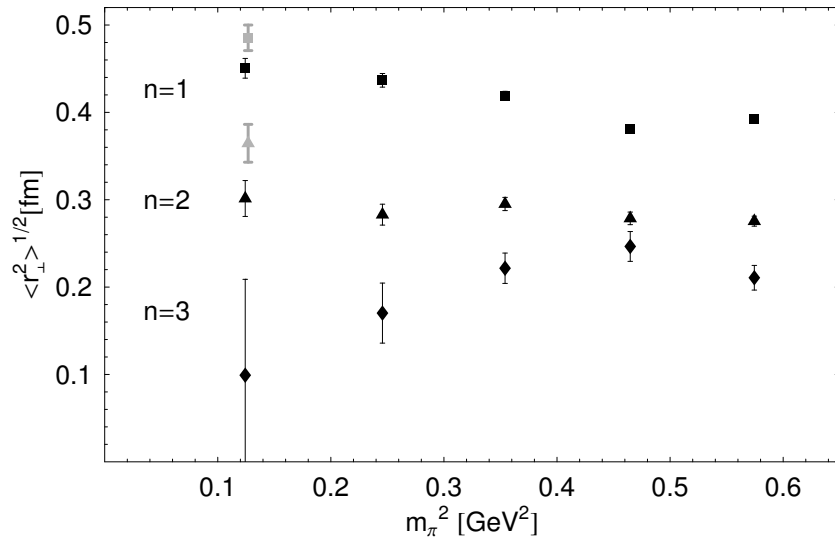
- simultaneous fits to q^2 and m_π dependence



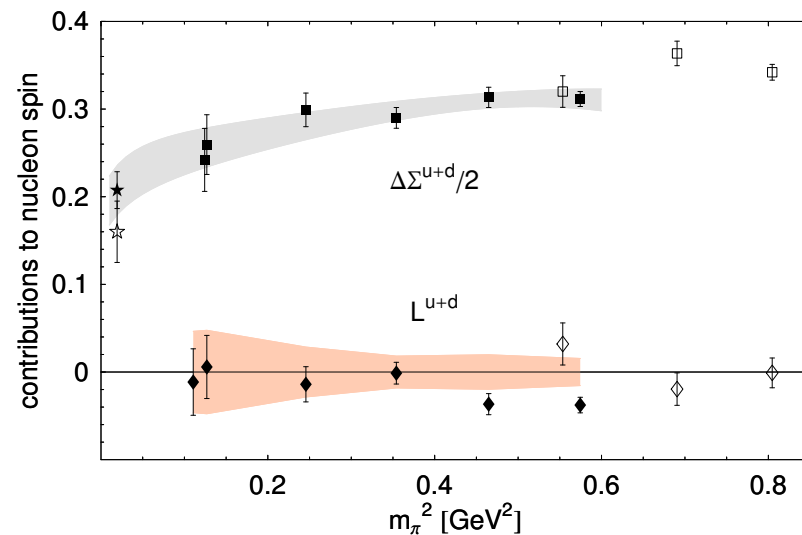
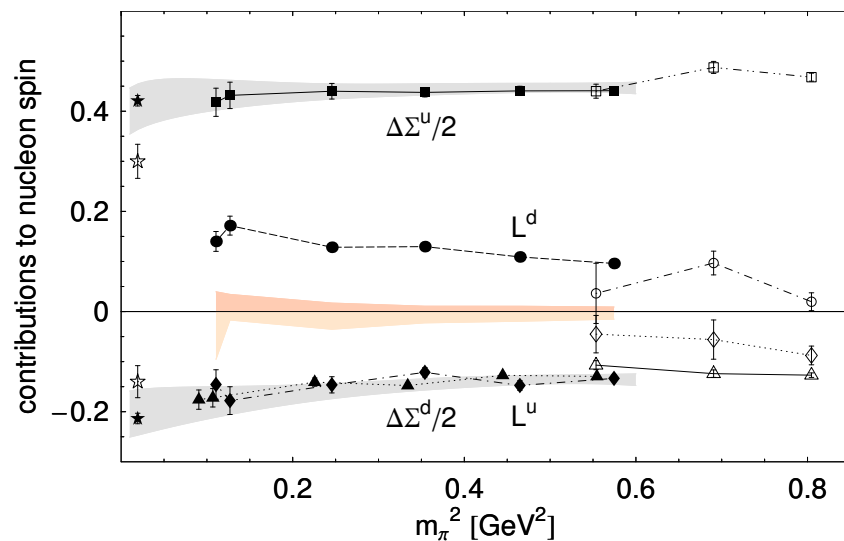
- comparing heavy baryon formulation (left) and covariant (right)

Transverse Structure and Spin Decomposition

- transverse size shows a significant dependence on n and hence x



- spin decomposition agrees with the latest HERMES results



Conclusions

- fits including oscillating exponentials provide a successful description of the oscillations observed in domain wall correlators
- replacing $\overset{\circ}{f}$, $\overset{\circ}{g}$, $\langle \overset{\circ}{x} \rangle$ with f , g , $\langle x \rangle$ in NLO terms of chipt give results that agree with experimental measurements of f_π , g_A , $\langle x \rangle_{u-d}$, $\langle x \rangle_{\Delta u - \Delta d}$
- the potential for predictions of $\langle \mathbf{1} \rangle_{\delta u - \delta d}$ and $\langle x \rangle_{\delta u - \delta d}$ is highlighted
- the relationship between m_π and q^2 dependence of generalized form factors is studied in both heavy baryon and convariant chipt
- the nucleon's transverse size shows a significant n (hence x) dependence
- the nucleon's spin decomposition agree with the latest HERMES results